

KAUST

CS 247 – Scientific Visualization Lecture 13: Scalar Field Visualization, Pt. 7

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Reading Assignment #7 (until Mar 19)



Read (required):

 Real-Time Volume Graphics, Chapter 1 (*Theoretical Background and Basic Approaches*), from beginning to 1.4.4 (inclusive)

Read (optional):

• Paper:

Nelson Max, Optical Models for Direct Volume Rendering, IEEE Transactions on Visualization and Computer Graphics, 1995 http://dx.doi.org/10.1109/2945.468400

Gradients as Differential Forms (1-Forms)

Interlude: Tensor Calculus

Tensors are multi-linear functions transforming contra- or covariantly

- *Contravariant* first-order tensor (*vector*) ۲
- *Covariant* first-order tensor (*covector*) ٠

The gradient vector is a contravariant vector $df = \frac{\partial f}{\partial x^i} dx^i$ The gradient 1-form is a covariant vector (a covector)

Very powerful; necessary for non-Cartesian coordinate systems / grids On (intrinsically) curved manifolds (sphere, ...): Cartesian coordinates not even possible

$\mathbf{v} = v^i \partial_i$

$$\mathbf{v} = v^i \, \mathbf{e}_i$$
$$\mathbf{w} = v_i \, \boldsymbol{\omega}^i$$

$$\mathbf{v} = v^i \mathbf{e}_i$$

Interlude: Tensor Calculus



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The gradient vector is a contravariant vector The gradient 1-form is a covariant vector (a covector) $\mathbf{v} = \overbrace{\frac{\partial f}{\partial x^i}}^{i} dx^i$

 $\mathbf{v} = v_i \mathbf{v}_i$ indexes! not exponents $\boldsymbol{\omega} = v_i \mathbf{\omega}^i$

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This is also the fundamental reason why in graphics a normal vector transforms differently: as a covector, not as a vector!

(typical graphics rule: **n** transforms with transpose of inverse matrix)

The Gradient as a Differential Form



The gradient as a *differential* (differential 1-form) is the "primary" concept (also "total differential" or "total derivative")

$$df = \frac{\partial f}{\partial x} \, dx + \frac{\partial f}{\partial y} \, dy + \frac{\partial f}{\partial z} \, dz$$

A differential 1-form is a scalar-valued linear function that takes a (direction) vector as input, and gives a scalar as output

Each of the 1-forms df, dx, dy, dz takes direction vector as input, gives scalar output

In the expression of the gradient df above, all 1-forms on the right-hand side get the same vector as input

df is simply a linear combination of the coordinate differentials dx, dy, dz

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The directional derivative and the gradient vector

$$D_{\mathbf{u}}f = df(\mathbf{u})$$
$$df(\mathbf{u}) = \nabla f \cdot \mathbf{u}$$

The gradient vector is then *defined*, such that:

$$\nabla f \cdot \mathbf{u} := df(\mathbf{u})$$





the function here is $f(x,y) = x^2 + y^2$ $\nabla f(x,y) = 2x\mathbf{i} + 2y\mathbf{j}$





the function here is $f(x,y) = x^2 + y^2$ $\nabla f(x,y) = 2x \mathbf{e}_x + 2y \mathbf{e}_y$ df(x,y) = 2x dx + 2y dy

 $df(r, \theta) = 2rdr + 0d\theta = 2rdr$





how about in polar coordinates?



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how about in polar coordinates?



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different 1-forms evaluated in some direction





 $\nabla f(r,\theta) = 2r\mathbf{e}_r + 0\frac{1}{r^2}\mathbf{e}_\theta = 2r\mathbf{e}_r$ $df(r,\theta) = 2rdr + 0d\theta = 2rdr$



Symmetric, covariant second-order tensor field: *defines* inner product on manifold (in each tangent space)

 $\langle \mathbf{v}, \mathbf{w} \rangle := \mathbf{g}(\mathbf{v}, \mathbf{w}) = \|\mathbf{v}\| \|\mathbf{w}\| \cos \theta$



Symmetric, covariant second-order tensor field: *defines* inner product on manifold (in each tangent space)

$$\langle \mathbf{v}, \mathbf{w} \rangle := \mathbf{g}(\mathbf{v}, \mathbf{w}) = \|\mathbf{v}\| \|\mathbf{w}\| \cos \theta$$

$$g_{ij} = \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix}$$

$$\|\mathbf{v}\|^2 = \langle \mathbf{v}, \mathbf{v} \rangle$$

$$= \mathbf{g}(\mathbf{v}, \mathbf{v})$$

$$= g_{ij} v^i v^j$$

$$= \mathbf{v}^T \mathbf{g} \mathbf{v}$$

$$(2D)$$



Symmetric, covariant second-order tensor field: *defines* inner product on manifold (in each tangent space)

$$\langle \mathbf{v}, \mathbf{w} \rangle := \mathbf{g}(\mathbf{v}, \mathbf{w}) = \|\mathbf{v}\| \|\mathbf{w}\| \cos \theta$$

$$g_{ij} = \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix} \begin{bmatrix} v^1 \\ v^2 \end{bmatrix}$$

$$\|\mathbf{v}\|^2 = \begin{bmatrix} v^1 & v^2 \end{bmatrix} \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix} \begin{bmatrix} v^1 \\ v^2 \end{bmatrix}$$

$$= \mathbf{g}(\mathbf{v}, \mathbf{v})$$

$$= g_{ij}v^iv^j$$

$$= \mathbf{v}^T \mathbf{g} \mathbf{v}$$

$$\|\mathbf{v}\|^2 = \begin{bmatrix} v^1 & v^2 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

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$$\|\mathbf{v}\|^2 = \begin{bmatrix} v^1 & v^2 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} v^1 \\ v^2 \end{bmatrix} = \mathbf{v}^T \mathbf{v}$$



Components of metric referred to coordinates

$$g_{ij} := \langle \mathbf{e}_i, \mathbf{e}_j \rangle$$

A second-order tensor field is bi-linear, i.e., linear in each (vector/covector) argument separately

From bi-linearity we immediately get:

$$\langle \mathbf{v}, \mathbf{w} \rangle = \mathbf{g} \left(v^i \mathbf{e}_i, w^j \mathbf{e}_j \right)$$

= $v^i w^j \langle \mathbf{e}_i, \mathbf{e}_j \rangle$
= $g_{ij} v^i w^j$

Gradient Vector from Differential 1-Form



The metric (and inverse metric) *lower* or *raise* indices (i.e., convert between covariant and contravariant tensors)

$$v^i = g^{ij} v_j$$
$$v_i = g_{ij} v^j$$

$$v^{i}\mathbf{e}_{i} = g^{ij}v_{j}\mathbf{e}_{i}$$
$$v_{i}\boldsymbol{\omega}^{i} = g_{ij}v^{j}\boldsymbol{\omega}^{i}$$

Inverse metric (contravariant)

$$[g^{ij}] = [g_{ij}]^{-1}$$

$$g^{ik}g_{kj}=\delta^i_j$$

Kronecker delta behaves like identity matrix

Gradient Vector from Differential 1-Form



So the gradient vector is

$$\nabla f = \left(g^{ij}\frac{\partial f}{\partial x^j}\right)\mathbf{e}_i$$

$$d\mathbf{r} = dx^i \,\mathbf{e}_i$$
$$d\mathbf{r}(\cdot) = dx^i(\cdot) \,\mathbf{e}_i$$

Directional derivative via inner product:

$$\langle \nabla f, \cdot \rangle = g_{kj} g^{ik} \frac{\partial f}{\partial x^i} dx^j(\cdot) \qquad \nabla f \cdot d\mathbf{r} = g_{kj} g$$

$$= \delta^i_j \frac{\partial f}{\partial x^i} dx^j(\cdot) \qquad = \delta^i_j \frac{\partial}{\partial x^i} dx^j(\cdot) \qquad = \delta^i_j \frac{\partial}{\partial x^i} dx^j(\cdot) \qquad = \frac{\partial}{\partial x^i} dx^i(\cdot) \qquad =$$

$$f \cdot d\mathbf{r} = g_{kj}g^{ik} \frac{\partial f}{\partial x^i} dx^j$$

= $\delta^i_j \frac{\partial f}{\partial x^i} dx^j$
= $\frac{\partial f}{\partial x^i} dx^i$

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Example: Polar Coordinates



Metric tensor and inverse metric for polar coordinates

$$g_{ij} = \begin{bmatrix} 1 & 0 \\ 0 & r^2 \end{bmatrix} \qquad \qquad g^{ij} = \begin{bmatrix} 1 & 0 \\ 0 & \frac{1}{r^2} \end{bmatrix}$$

Gradient vector from 1-form: raise index with inverse metric

$$\nabla f = \begin{bmatrix} \frac{\partial f}{\partial r} & \frac{\partial f}{\partial \theta} \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & \frac{1}{r^2} \end{bmatrix} \begin{bmatrix} \mathbf{e}_r \\ \mathbf{e}_{\theta} \end{bmatrix} = \frac{\partial f}{\partial r} \mathbf{e}_r + \frac{1}{r^2} \frac{\partial f}{\partial \theta} \mathbf{e}_{\theta}$$
 not normalized!

Example: Polar Coordinates



Metric tensor and inverse metric for polar coordinates

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Gradient vector from 1-form: raise index with inverse metric

$$\nabla f(r,\theta) = \frac{\partial f(r,\theta)}{\partial r} \mathbf{e}_r(r,\theta) + \frac{1}{r^2} \frac{\partial f(r,\theta)}{\partial \theta} \mathbf{e}_\theta(r,\theta)$$

and don't forget that all of this is position-dependent!

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Tensor Calculus

Highly recommended:

Very nice book,

complete lecture on Youtube!

Pavel Grinfeld

Introduction to Tensor Analysis and the Calculus of Moving Surfaces

D Springer



Thank you.

Thanks for material

- Helwig Hauser
- Eduard Gröller
- Daniel Weiskopf
- Torsten Möller
- Ronny Peikert
- Philipp Muigg
- Christof Rezk-Salama