

KAUST

CS 247 – Scientific Visualization Lecture 28: Vector / Flow Visualization, Pt. 7

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Reading Assignment #15++ (1)



Reading suggestions:

- Data Visualization book, Chapter 6.7
- J. van Wijk: *Image-Based Flow Visualization*, ACM SIGGRAPH 2002

http://www.win.tue.nl/~vanwijk/ibfv/ibfv.pdf

• T. Günther, A. Horvath, W. Bresky, J. Daniels, S. A. Buehler: Lagrangian Coherent Structures and Vortex Formation in High Spatiotemporal-Resolution Satellite Winds of an Atmospheric Karman Vortex Street, 2021

https://www.essoar.org/doi/10.1002/essoar.10506682.2

 H. Bhatia, G. Norgard, V. Pascucci, P.-T. Bremer: *The Helmholtz-Hodge Decomposition – A Survey*, TVCG 19(8), 2013

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https://doi.org/10.1109/TVCG.2012.316
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• Work through online tutorials of multi-variable partial derivatives, grad, div, curl, Laplacian:

https://www.khanacademy.org/math/multivariable-calculus/multivariable-derivatives https://www.youtube.com/watch?v=rB83DpBJQsE(3Blue1Brown)

• Matrix exponentials:

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https://www.youtube.com/watch?v=0850WBJ2ayo (3Blue1Brown)
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Reading Assignment #15++ (2)



Reading suggestions:

 Tobias Günther, Irene Baeza Rojo: *Introduction to Vector Field Topology* https://cgl.ethz.ch/Downloads/Publications/Papers/2020/Gun20b/Gun20b.pdf

 Roxana Bujack, Lin Yan, Ingrid Hotz, Christoph Garth, Bei Wang: State of the Art in Time-Dependent Flow Topology: Interpreting Physical Meaningfulness Through Mathematical Properties https://onlinelibrary.wiley.com/doi/epdf/10.1111/cgf.14037

- B. Jobard, G. Erlebacher, M. Y. Hussaini: Lagrangian-Eulerian Advection of Noise and Dye Textures for Unsteady Flow Visualization http://dx.doi.org/10.1109/TVCG.2002.1021575
- Anna Vilanova, S. Zhang, Gordon Kindlmann, David Laidlaw: An Introduction to Visualization of Diffusion Tensor Imaging and Its Applications http://vis.cs.brown.edu/docs/pdf/Vilanova-2005-IVD.pdf

Lagrangian vs. Eulerian

Lagrangian vs. Eulerian

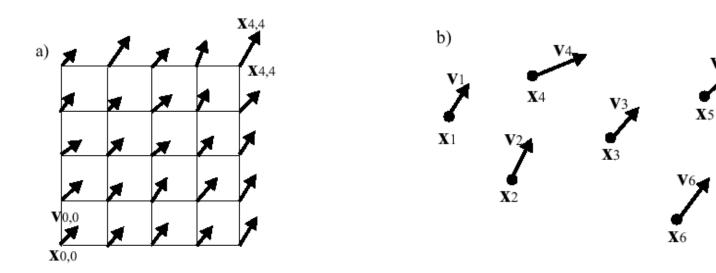


Eulerian

- Flow properties given at fixed spatial positions (grid points)
- Partial time derivative

Lagrangian

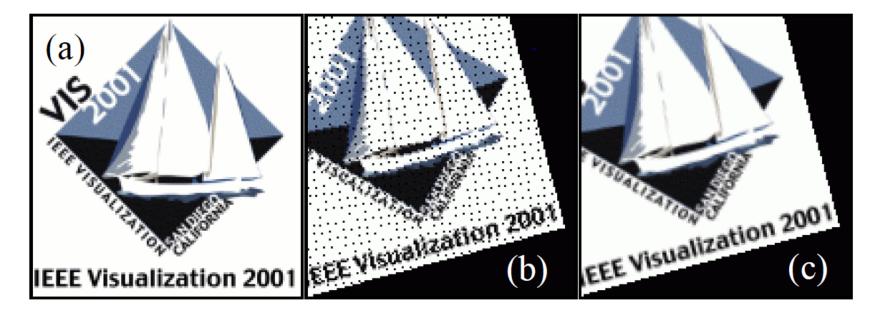
- Flow properties given for each particle (particles are moving)
- Material time derivative



Lagrangian vs. Eulerian



- Lagrangian: move along with the particle
- Eulerian: consider fixed point in space, look at particles moving through



 Example for pixels: rotate image (a), Lagrangian: move pixels forward (b), Eulerian: fetch pixels from backward direction (c)



The material time derivative (convective derivative) gives the rate of change when following a particle in the flow

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + (\mathbf{u} \cdot \nabla)$$



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$$\frac{D}{Dt} = \frac{\partial}{\partial t} + (\mathbf{u} \cdot \nabla)$$
$$\frac{DT}{Dt} = \frac{\partial T}{\partial t} + (\mathbf{u} \cdot \nabla) T$$



The material time derivative (convective derivative) gives the rate of change when following a particle in the flow

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + (\mathbf{u} \cdot \nabla)$$
$$\frac{DT}{Dt} = \frac{\partial T}{\partial t} + (\mathbf{u} \cdot \nabla) T$$
$$\frac{DT}{Dt} = \frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z}$$



$$dT = \frac{\partial T}{\partial t}dt + \frac{\partial T}{\partial x}dx + \frac{\partial T}{\partial y}dy + \frac{\partial T}{\partial z}dz$$



$$dT = \frac{\partial T}{\partial t}dt + \frac{\partial T}{\partial x}dx + \frac{\partial T}{\partial y}dy + \frac{\partial T}{\partial z}dz$$
$$dT = \frac{\partial T}{\partial t}dt + \frac{\partial T}{\partial x}\frac{dx}{dt}dt + \frac{\partial T}{\partial y}\frac{dy}{dt}dt + \frac{\partial T}{\partial z}\frac{dz}{dt}dt$$



$$dT = \frac{\partial T}{\partial t}dt + \frac{\partial T}{\partial x}dx + \frac{\partial T}{\partial y}dy + \frac{\partial T}{\partial z}dz$$
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$$\frac{\partial T}{\partial t} = \frac{\partial T}{\partial t} + \frac{\partial T}{\partial x}\frac{dx}{dt} + \frac{\partial T}{\partial y}\frac{dy}{dt} + \frac{\partial T}{\partial z}\frac{dz}{dt}$$



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$$\frac{dT}{dt} = \frac{\partial T}{\partial t} + \frac{\partial T}{\partial x}\frac{dx}{dt} + \frac{\partial T}{\partial y}\frac{dy}{dt} + \frac{\partial T}{\partial z}\frac{dz}{dt}$$
$$u := \frac{dx}{dt}, \quad v := \frac{dy}{dt}, \quad w := \frac{dz}{dt}$$



Actually, nothing else than application of the multi-variable chain rule:

We are given T(x, y, z, t) with four independent variables;

But now we want to go along a parameterized path with parameter t, so x, y, z become dependent variables: x(t), y(t), z(t)

Along this path, our goal is now to compute the derivative of the function

T(x(t), y(t), z(t), t) with t as only independent variable:

$$\begin{aligned} \frac{d}{dt}T\left(x(t), y(t), z(t), t\right) &= \\ \frac{\partial}{\partial t}T(x, y, z, t) + \frac{\partial}{\partial x}T(x, y, z, t)\frac{d}{dt}x(t) + \frac{\partial}{\partial y}T(x, y, z, t)\frac{d}{dt}y(t) + \frac{\partial}{\partial z}T(x, y, z, t)\frac{d}{dt}z(t) \\ u(t) &:= \frac{dx(t)}{dt}, \quad v(t) := \frac{dy(t)}{dt}, \quad w(t) := \frac{dz(t)}{dt} \end{aligned}$$

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Advection



Advection equation; velocity field **u**(*x*, *y*, *z*, *t*), no change following particle, just advection: set material derivative = 0:

$$\frac{\partial T}{\partial t} + \left(\mathbf{u} \cdot \nabla\right) T = 0$$

In the Navier-Stokes equations: "self-advection" of velocity

• Advect scalar components of velocity field individually (actually two equations in 2D, three equations in 3D)

$$\frac{\partial \mathbf{u}}{\partial t} = -\big(\mathbf{u} \cdot \nabla\big)\mathbf{u}$$

this is equivalent to saying that the acceleration is zero!

Fluid Simulation Basics

Fluid Simulation and Rendering

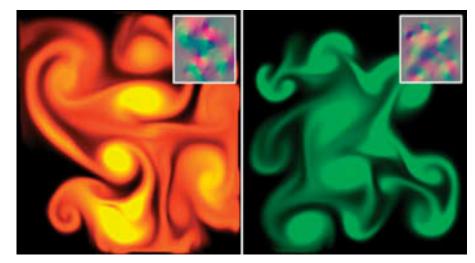


Compute advection of fluid

- (Incompressible or compressible) Navier-Stokes solvers
- Lattice Boltzmann Method (LBM)

Discretized domain

- Velocity, pressure
- Dye, smoke density, vorticity, ...



Courtesy Mark Harris

Fluid Simulation: Navier Stokes (1)



Incompressible (divergence-free) Navier Stokes equations

$$\frac{\partial \mathbf{u}}{\partial t} = -(\mathbf{u} \cdot \nabla)\mathbf{u} - \frac{1}{\rho}\nabla\rho + \nu\nabla^2\mathbf{u} + \mathbf{F},$$
$$\nabla \cdot \mathbf{u} = 0,$$

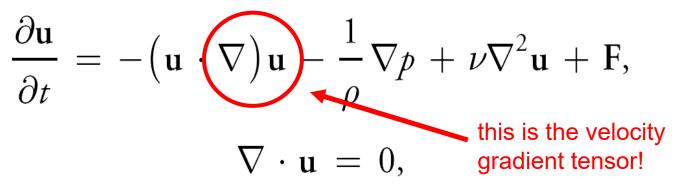
Components:

- Self-advection of velocity (i.e., advection of velocity according to velocity)
- Pressure gradient (force due to pressure differences)
- Diffusion of velocity due to viscosity (for viscous fluids, i.e., not inviscid)
- Application of (arbitrary) external forces, e.g., gravity, user input, etc.

Fluid Simulation: Navier Stokes (1)



Incompressible (divergence-free) Navier Stokes equations



Components:

- Self-advection of velocity (i.e., advection of velocity according to velocity)
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Fluid Simulation: Navier Stokes (2)



Given a (Cartesian) coordinate system, the momentum equation can be seen as a system of equations (2 equations in 2D, 3 equations in 3D)

For 2D (Cartesian):

$$\frac{\partial u}{\partial t} = -(\mathbf{u} \cdot \nabla) u - \frac{1}{\rho} (\nabla p)_x + \nu \nabla^2 u + f_x,$$

$$\frac{\partial v}{\partial t} = -(\mathbf{u} \cdot \nabla) v - \frac{1}{\rho} (\nabla p)_{y} + \nu \nabla^{2} v + f_{y}.$$

these are PDEs!

Vector Fields, Vector Calculus, and Dynamical Systems

Some Vector Calculus (1)

Gradient (scalar field \rightarrow vector field)

- Direction of steepest ascent; magnitude = rate
- Conservative vector field: gradient of some scalar (potential) function

Divergence (vector field \rightarrow scalar field)

- Volume density of outward flux: "exit rate: source? sink?"
- Incompressible/solenoidal/divergence-free vector field: div u = 0 can express as curl (next slide) of some vector (potential) function

Laplacian (scalar field \rightarrow scalar field)

- Divergence of gradient
- Measure for difference between point and its neighborhood

$$\nabla \cdot \mathbf{u} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}$$

$$\nabla^2 p = \frac{\partial^2 p}{\partial x^2} + \frac{\partial^2 p}{\partial y^2}$$

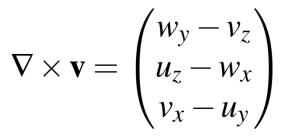
 $\nabla p = \left(\frac{\partial p}{\partial x}, \ \frac{\partial p}{\partial y}\right)$

Some Vector Calculus (2)



Curl (vector field \rightarrow vector field)

- Circulation density at a point (*vorticity*)
- If curl vanishes everywhere: irrotational/curl-free field
- Every conservative (path-independent) field is irrotational (and vice versa if domain is simply connected)



these are partial derivatives!

st re $-\frac{4}{4}$ $-\frac{2}{2}$ 0 2 4

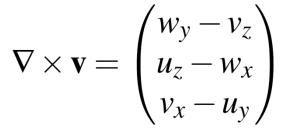
Example: curl = const everywhere

Some Vector Calculus (3)



Curl (vector field \rightarrow vector field)

- Circulation density at a point (*vorticity*)
- If curl vanishes everywhere: irrotational/curl-free field
- Every conservative (path-independent) field is irrotational (and vice versa if domain is simply connected)



these are partial derivatives!

Example: curl not always "obviously rotational"

Some Vector Calculus (4)



Curl (vector field \rightarrow vector field)

- Circulation density at a point (vorticity)
- If curl vanishes everywhere: irrotational/curl-free field
- Every conservative (path-independent) field is irrotational (and vice versa if domain is simply connected)

Example: non-obvious curl-free field

[this domain is **not** simply connected! it is the "punctured plane", i.e., the point (0,0) is not in the domain]

$$7 \times \mathbf{v} = \begin{pmatrix} w_y - v_z \\ u_z - w_x \\ v_x - u_y \end{pmatrix}$$

these are partial derivatives!

$$\mathbf{V}(x,y,z) = \frac{(-y,x,0)}{x^2 + y^2}$$

not defined at (x,y) = (0,0)

$$v_x = u_y \qquad \nabla \times \mathbf{v} = \mathbf{0}$$

velocity gradient ∇v is symmetric (see later)

Some Vector Calculus (5)

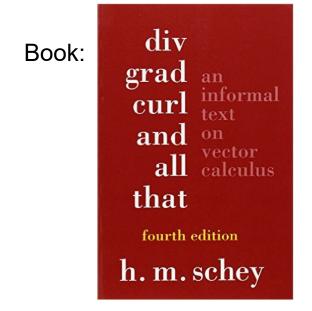


Curl (vector field \rightarrow vector field)

- Circulation density at a point (*vorticity*)
- If curl vanishes everywhere: irrotational/curl-free field
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 $\nabla \times \mathbf{v} = \begin{pmatrix} w_y - v_z \\ u_z - w_x \\ v_x - u_y \end{pmatrix}$

these are partial derivatives!



Interactive tutorial on curl: http://mathinsight.org/curl_idea

Fundamental theorem of vector calculus: Helmholtz decomposition: Any vector field can be expressed as the sum of a solenoidal (*divergence-free*) vector field and an irrotational (*curl-free*) vector field (Helmholtz-Hodge: plus *harmonic* vector field)

Vector Fields and Dynamical Systems (1)



Velocity gradient tensor, (vector field \rightarrow tensor field)

- Gradient of vector field: how does the vector field change?
- In Cartesian coordinates: *spatial partial derivatives (Jacobian matrix)*

$$\nabla \mathbf{v} (x, y, z) = \begin{pmatrix} u_x & u_y & u_z \\ v_x & v_y & v_z \\ w_x & w_y & w_z \end{pmatrix}$$
these are partial derivatives!

• Can be decomposed into symmetric part + anti-symmetric part

 $\nabla \mathbf{v} = \mathbf{D} + \mathbf{S}$

velocity gradient tensor

sym.: $\mathbf{D} = \frac{1}{2} (\nabla \mathbf{v} + (\nabla \mathbf{v})^T)$ skew-sym.: $S = \frac{1}{2} (\nabla v - (\nabla v)^T)$ rotation: *vorticity/spin tensor*

deform : rate-of-strain tensor

Vector Fields and Dynamical Systems (2)



thoos or

Vorticity/spin/angular velocity tensor

- Antisymmetric part of velocity gradient tensor
- Corresponds to vorticity/curl/angular velocity (beware of factor 1/2)

$$\mathbf{S} = \frac{1}{2} \left(\nabla \mathbf{V} - (\nabla \mathbf{V})^{\mathrm{T}} \right)$$
partial
derivatives!

$$\mathbf{S} = \frac{1}{2} \begin{pmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{pmatrix} \quad \boldsymbol{\omega} = \begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix} = \nabla \times \mathbf{v} = \begin{pmatrix} w_y - v_z \\ u_z - w_x \\ v_x - u_y \end{pmatrix}$$

S acts on vector like cross product with ω : S • = $\frac{1}{2}\omega \times$

$$\mathbf{v}^{(r)} = \mathbf{S} \cdot d\mathbf{r} = \frac{1}{2} \left[\nabla \mathbf{v} - (\nabla \mathbf{v})^T \right] \cdot d\mathbf{r} = \frac{1}{2} \boldsymbol{\omega} \times d\mathbf{r}$$

Angular Velocity of Rigid Body Rotation

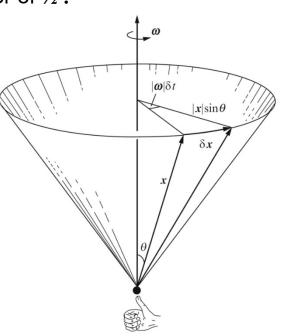
Rate of rotation

- Scalar ω: angular displacement per unit time (rad s⁻¹)
 - Angle Θ at time t is $\Theta(t) = \omega t$; $\omega = 2\pi f$ where f is the frequency (f = 1/T; s⁻¹)
- Vector $\boldsymbol{\omega}$: axis of rotation; magnitude is angular speed (if $\boldsymbol{\omega}$ is curl: speed x2)
 - Beware of different conventions that differ by a factor of $\frac{1}{2}$!

Cross product of $\frac{1}{2}\omega$ with vector to center of rotation (r) gives linear velocity vector v (tangent)

X X X X A A A

$$\mathbf{v}^{(r)} = rac{1}{2} \, oldsymbol{\omega} \, imes d\mathbf{r}$$



Velocity Gradient Tensor and Components (1)



Velocity gradient tensor

(here: in Cartesian coordinates)

$$\nabla \mathbf{v} = \begin{bmatrix} \frac{\partial}{\partial x} v^{x} & \frac{\partial}{\partial y} v^{x} & \frac{\partial}{\partial z} v^{x} \\ \frac{\partial}{\partial x} v^{y} & \frac{\partial}{\partial y} v^{y} & \frac{\partial}{\partial z} v^{y} \\ \frac{\partial}{\partial x} v^{z} & \frac{\partial}{\partial y} v^{z} & \frac{\partial}{\partial z} v^{z} \end{bmatrix}$$

these are the same partial derivatives as before!

$$\nabla \mathbf{v} = \frac{1}{2} \left(\nabla \mathbf{v} + (\nabla \mathbf{v})^T \right) + \frac{1}{2} \left(\nabla \mathbf{v} - (\nabla \mathbf{v})^T \right)$$

Velocity Gradient Tensor and Components (2)

Rate-of-strain (rate-of-deformation) tensor

(symmetric part; here: in Cartesian coordinates)

$$\mathbf{D} = \frac{1}{2} \begin{bmatrix} 2\frac{\partial}{\partial x}v^{x} & \frac{\partial}{\partial y}v^{x} + \frac{\partial}{\partial x}v^{y} & \frac{\partial}{\partial z}v^{x} + \frac{\partial}{\partial x}v^{z} \\ \frac{\partial}{\partial x}v^{y} + \frac{\partial}{\partial y}v^{x} & 2\frac{\partial}{\partial y}v^{y} & \frac{\partial}{\partial z}v^{y} + \frac{\partial}{\partial y}v^{z} \\ \frac{\partial}{\partial x}v^{z} + \frac{\partial}{\partial z}v^{x} & \frac{\partial}{\partial y}v^{z} + \frac{\partial}{\partial z}v^{y} & 2\frac{\partial}{\partial z}v^{z} \end{bmatrix}$$

$$tr(\mathbf{D}) = \nabla \cdot \mathbf{v}$$

Velocity Gradient Tensor and Components (3)



Vorticity tensor (spin tensor)

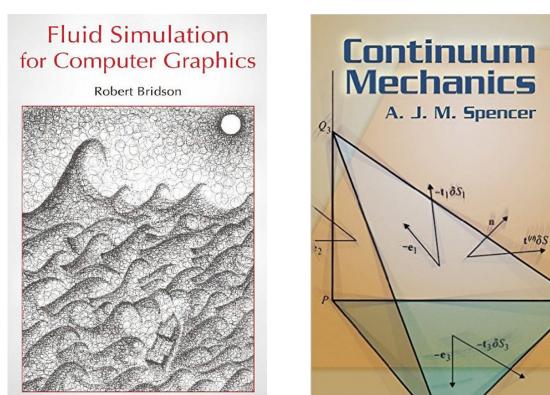
(skew-symmetric part; here: in Cartesian coordinates)

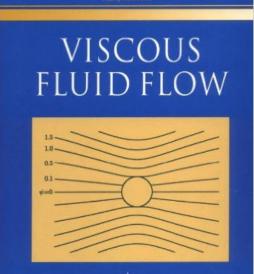
$$\mathbf{S} = \frac{1}{2} \begin{bmatrix} 0 & \frac{\partial}{\partial y} v^{x} - \frac{\partial}{\partial x} v^{y} & \frac{\partial}{\partial z} v^{x} - \frac{\partial}{\partial x} v^{z} \\ \frac{\partial}{\partial x} v^{y} - \frac{\partial}{\partial y} v^{x} & 0 & \frac{\partial}{\partial z} v^{y} - \frac{\partial}{\partial y} v^{z} \\ \frac{\partial}{\partial x} v^{z} - \frac{\partial}{\partial z} v^{x} & \frac{\partial}{\partial y} v^{z} - \frac{\partial}{\partial z} v^{y} & 0 \end{bmatrix}$$

$$\mathbf{S} = \frac{1}{2} \begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix} \qquad \boldsymbol{\omega} \equiv \nabla \times \mathbf{v}$$

Recommended Books (1)



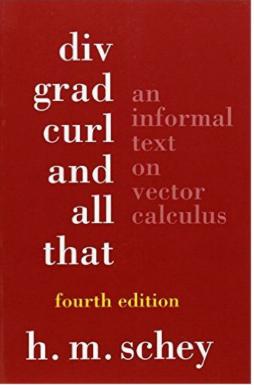


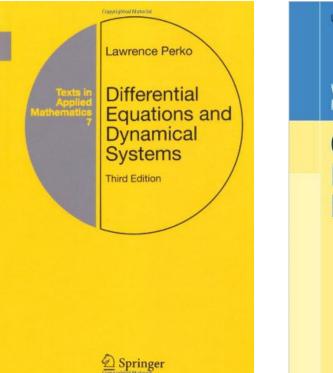


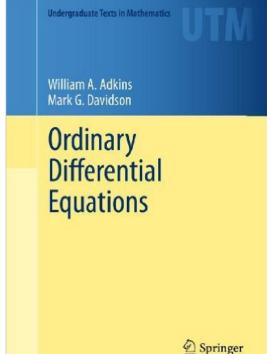
Tasos C. Papanastasiou Georgios C. Georgiou Andreas N. Alexandrou

Recommended Books (2)









Thank you.

Thanks for material

- Helwig Hauser
- Eduard Gröller
- Daniel Weiskopf
- Torsten Möller
- Ronny Peikert
- Philipp Muigg
- Christof Rezk-Salama